

In this lesson, we will:

- Derive the **vorticity equation** in tensor notation
- Review some **vector identities** and the **epsilon-delta relation**

Derivation of the Vorticity Equation in Tensor Notation

Start with the incompressible Navier-Stokes equation with constant properties:

$$\rho \frac{\partial u_i}{\partial t} + \rho u_j \frac{\partial u_i}{\partial x_j} = -\frac{\partial p}{\partial x_i} + \rho g_i + \mu \frac{\partial^2 u_i}{\partial x_j \partial x_j}$$



$$u_j \frac{\partial u_i}{\partial x_j} = \frac{\partial}{\partial x_i} \left(\frac{1}{2} u_j u_j \right) - u_j \epsilon_{ijn} \omega_n$$

$$\rho g_i = -\rho \frac{\partial}{\partial x_i} (gz)$$

→ move to RHS

$$\rho \frac{\partial u_i}{\partial t} + \rho \frac{\partial}{\partial x_i} \left(\frac{1}{2} u_j u_j \right) = -\frac{\partial p}{\partial x_i} - \rho \frac{\partial}{\partial x_i} (gz) + \rho u_j \epsilon_{ijn} \omega_n + \mu \frac{\partial^2 u_i}{\partial x_j \partial x_j} \quad (1)$$

TAKE CURL OF EQ.(1) $\nabla \times [Eq.(1)]$

$$\left(\frac{\partial}{\partial x_m} \frac{\partial}{\partial t} = \frac{\partial}{\partial t} \frac{\partial}{\partial x_m} \right)$$

$$\epsilon_{kmi} \frac{\partial}{\partial x_m} [Eq.(1)]$$

⇒ VECTOR EQ WITH FREE INDEX K

$$\rho \left[\begin{aligned} & \textcircled{1} \epsilon_{kmi} \frac{\partial}{\partial t} \frac{\partial u_i}{\partial x_m} + \textcircled{2} \epsilon_{kmi} \frac{\partial}{\partial x_m} \frac{\partial}{\partial x_i} \left(\frac{1}{2} u_j u_j \right) - \textcircled{3} \epsilon_{kmi} \frac{\partial}{\partial x_m} \left(u_j \epsilon_{ijn} \omega_n \right) \\ & = -\textcircled{4} \epsilon_{kmi} \frac{\partial}{\partial x_m} \frac{\partial}{\partial x_i} p - \textcircled{5} \rho \epsilon_{kmi} \frac{\partial^2 (gz)}{\partial x_m \partial x_i} + \mu \textcircled{6} \epsilon_{kmi} \frac{\partial^3 u_i}{\partial x_m \partial x_j \partial x_j} \end{aligned} \right] \quad (2)$$

• TERM ①

$$\rho \frac{\partial}{\partial t} \left(\epsilon_{kmi} \frac{\partial u_i}{\partial x_m} \right) \rightarrow \omega_k$$

$$\textcircled{1} = \rho \frac{\partial \omega_k}{\partial t}$$

• TERM ②

$$\rho \left(\epsilon_{kmi} \frac{\partial}{\partial x_m} \frac{\partial}{\partial x_i} \left(\frac{1}{2} u_j u_j \right) \right)$$

ANTISYMMETRIC IN $m \ \& \ i$

SYMMETRIC IN $m \ \& \ i$

RECALL, THE CONTRACTED PRODUCT OF A SYMMETRIC & ANTISYMMETRIC TENSOR IS ZERO

$$\therefore \textcircled{2} = 0$$

• TERM ④

→ SAME ARGUMENT AS FOR TERM ② →

$$\textcircled{4} = 0$$

• TERM ⑤

.....

$$\textcircled{5} = 0$$

• TERM ⑥

$$\mu \epsilon_{kmi} \frac{\partial^3}{\partial x_m \partial x_j \partial x_j} u_i = \mu \frac{\partial}{\partial x_j} \frac{\partial}{\partial x_j} \left(\epsilon_{kmi} \frac{\partial u_i}{\partial x_m} \right)$$

ω_k !

$$\textcircled{6} = \mu \frac{\partial^2 \omega_k}{\partial x_j \partial x_j}$$

• TERM ③

$$-\rho \epsilon_{kmi} \frac{\partial}{\partial x_m} (u_j \epsilon_{ijn} \omega_n) = -\rho \epsilon_{kmi} \epsilon_{ijn} \frac{\partial}{\partial x_m} (u_j \omega_n)$$

Recall the ϵ - δ relation

$$\epsilon_{ijk} \epsilon_{klm} = \delta_{il} \delta_{jm} - \delta_{im} \delta_{jl}$$

⇓

$$\epsilon_{kmi} \epsilon_{ijn} = \delta_{kj} \delta_{mn} - \delta_{kn} \delta_{mj}$$

CHANGE ALL i TO k

CHANGE ALL j TO m

⇓

$$\textcircled{3} = -\rho \left[\delta_{kj} \delta_{mn} \frac{\partial}{\partial x_m} (u_j \omega_n) - \delta_{kn} \delta_{mj} \frac{\partial}{\partial x_m} (u_j \omega_n) \right]$$

contracts

non-zero only if $j=k$

non-zero only if $j=m$

$$(3) = -\rho \left[\int_{mn} \frac{\partial}{\partial x_m} (u_k \omega_n) - \int_{kn} \frac{\partial}{\partial x_j} (u_j \omega_n) \right]$$

$$(3) = -\rho \left[\frac{\partial}{\partial x_n} (u_k \omega_n) - \frac{\partial}{\partial x_j} (u_j \omega_k) \right]$$

$$(3) = -\rho \left[\frac{\partial}{\partial x_j} (u_k \omega_j) - \frac{\partial}{\partial x_j} (u_j \omega_k) \right]$$

USE PRODUCT RULE ON BOTH TERMS

$$(3) = -\rho \left[\cancel{u_k \frac{\partial \omega_j}{\partial x_j}} + \omega_j \frac{\partial u_k}{\partial x_j} - \omega_k \frac{\partial u_j}{\partial x_j} - u_j \frac{\partial \omega_k}{\partial x_j} \right]$$

0 for incompressible flow

$$u_k \frac{\partial}{\partial x_j} \left(\epsilon_{imj} \frac{\partial u_m}{\partial x_i} \right) = 0$$

$$(3) = -\rho \omega_j \frac{\partial u_k}{\partial x_j} + \rho u_j \frac{\partial \omega_k}{\partial x_j}$$

(2) becomes $\rho \left[\frac{D\omega_k}{Dt} + u_j \frac{\partial \omega_k}{\partial x_j} \right] = \rho \omega_j \frac{\partial u_k}{\partial x_j} + \mu \frac{\partial^2 \omega_k}{\partial x_j \partial x_j}$

$\mu = \rho \nu$

$$\star \frac{D\omega_k}{Dt} = \omega_j \frac{\partial u_k}{\partial x_j} + \nu \frac{\partial^2 \omega_k}{\partial x_j \partial x_j}$$

$$\Rightarrow \frac{D\vec{\omega}}{Dt} = (\vec{\omega} \cdot \nabla) \vec{u} + \nu \nabla^2 \vec{\omega}$$

VORTICITY EQUATION FOR INCOMPRESSIBLE, NEWTONIAN FLOW
(with constant properties)