

Today, we will:

- Continue with the example problem – the Bénard problem – temperature instability between two infinite plates: Apply the method of normal modes and solve.

We begin with Eqs. (7) and (8) on the handout (semi-normalized eqs. for the disturbance variables w and T'):

$$\frac{\partial T'}{\partial t} - \frac{\Gamma d^2}{\kappa} w = \nabla^2 T' \quad (7), \text{ and } \left(\frac{1}{\text{Pr}} \frac{\partial}{\partial t} - \nabla^2 \right) \nabla^2 w = \frac{g \alpha d^2}{\nu} \nabla_H^2 T' \quad (8).$$

along with their BCs: $w = \frac{\partial w}{\partial z} = T' = 0$ at $z = \pm \frac{1}{2}$.

Apply the Method of Normal Modes → disturbances are periodic in x, y may be periodic and/or growing in t

Let $w = \hat{w}(z) e^{ikx + ily + \sigma t}$
 $T' = \hat{T}(z)$ " "

Where $\hat{w}(z), \hat{T}(z)$ are complex amplitudes

k, l are x, y components of wavenumber vector $\vec{k} = (k, l, m)$

Here k, l must be real → no growth in x, y directions

$\sigma = \text{growth rate} = \sigma_r + i\sigma_i$ in general (may be growing or decaying in time)
 [temporal stability analysis]

Plug these into Eqs (7) & (8)

$$(7) \rightarrow \frac{\partial T'}{\partial t} - \frac{\Gamma d^2}{\kappa} w = \nabla^2 T'$$

$$\hat{T}(z) e^{ikx + ily + \sigma t} - \frac{\Gamma d^2}{\kappa} \hat{w}(z) e^{ikx + ily + \sigma t} = \hat{T}(z) e^{ikx + ily + \sigma t} (ik)^2 + \dots$$

every term has $e^{ikx + ily + \sigma t}$ so we can cancel it out

$i^2 k^2 = -k^2$

also, define $K = \sqrt{k^2 + l^2}$ = magnitude of the wavenumber vector

We get $\left[\delta - \left(\frac{d^2}{dz^2} - K^2 \right) \right] \hat{T} = \frac{\Gamma d^2}{K} \hat{w}$ (9)

Let $D = \frac{d}{dz}$

$W = \frac{\Gamma d^2}{K} \hat{w}$

$Ra = \text{Rayleigh \#} = \frac{g \alpha |\Delta T| d^3}{K \nu} = \frac{g \alpha \Gamma d^4}{K \nu}$

since $\Gamma = \frac{\Delta T}{d}$

(9) becomes Eq. (11) on handout

i. (8) " " (12) " "

BC's become

$W = DW = \hat{T} = 0 @ z = \pm 1/2$

Comments: 1) No longer any distinction between x & y directions

$k = \text{wavenumber in x-direction}$
 $l = \text{" " " y-direction}$
 } combined into $K^2 = \sqrt{k^2 + l^2}$

(because nothing special about x or y directions)

2) x, y, & t have dropped out! z is the only remaining indep. variable

Now we have ODEs!

originally,

5 pdes
5 unknowns

⇒

2 pdes
2 unknowns

⇒

2 ODEs
2 unknowns

u, T', p as fncs of x, y, z, t

w, T' as fncs of x, y, z, t

$W(z)$ & $\hat{T}(z)$

3) It can be proven (see Kundu, pg. 460-461) that σ must be real

$$\sigma = \sigma_r + i\sigma_i = \sigma_r = \underline{\text{real}}$$

recall $w = \hat{w}(z) e^{ikx + iky + \sigma t}$

growth i. jumps in time + periodicity in time

sinusoidal waves in x & y directions

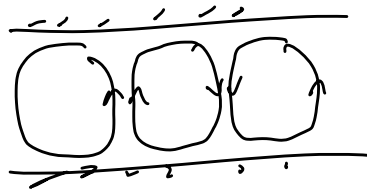
But $e^{\sigma t} = e^{\sigma_r t} + i\sigma_i t$

\therefore NO PERIODICITY IN TIME

\therefore we conclude that the instability will cause the basic state to jump to another steady state rather than an oscillating state

agrees with experiment \rightarrow get cells

We call this The principle of \star exchange of stabilities



Also, since $\sigma = \sigma_r = \text{real}$, σ_r determines the stability

$$e^{\sigma t} \rightarrow \begin{cases} \text{if } \sigma < 0 & \text{stable} \\ \text{if } \sigma > 0 & \text{unstable} \\ \text{if } \sigma = 0 & \text{neutrally or marginally stable} \end{cases}$$

\star Let's examine the marginal stability case

\star Let's set or force $\sigma = 0$ to find the marginal stability condition

Eqs (11) & (12) $\Rightarrow (D^2 - K^2) \hat{T} = -W$ (11) \star marginal

$(D^2 - K^2)^2 W = Ra K^2 \hat{T}$ (12) \star

To solve (11 m): (2 m): Take $(D^2 - K^2)$ of (12 m):

$$(D^2 - K^2)^3 W = R_a K^2 (D^2 - K^2) T = -W \text{ from Eq. (11 m)}$$

$$\therefore (D^2 - K^2)^3 W = -R_a K^2 W \rightarrow ! \text{ we've eliminated } T'$$

ONE ODE FOR $W(x)$

$$\left[(D^2 - K^2)^3 + R_a K^2 \right] W(x) = 0 \quad (13 m)$$

★ a linear, homogeneous o.d.e. with constant coefficients

Solve using the "auxiliary eq." for q

$$(q^2 - K^2)^3 + R_a K^2 = 0 \rightarrow \text{Find the roots of this eq.}$$

6th-order $\rightarrow \therefore$ get 6 roots $q_1, q_2, q_3, \dots, q_6$

$$\text{Solution is } W(x) = C_1 e^{q_1 x} + C_2 e^{q_2 x} + \dots + C_6 e^{q_6 x}$$

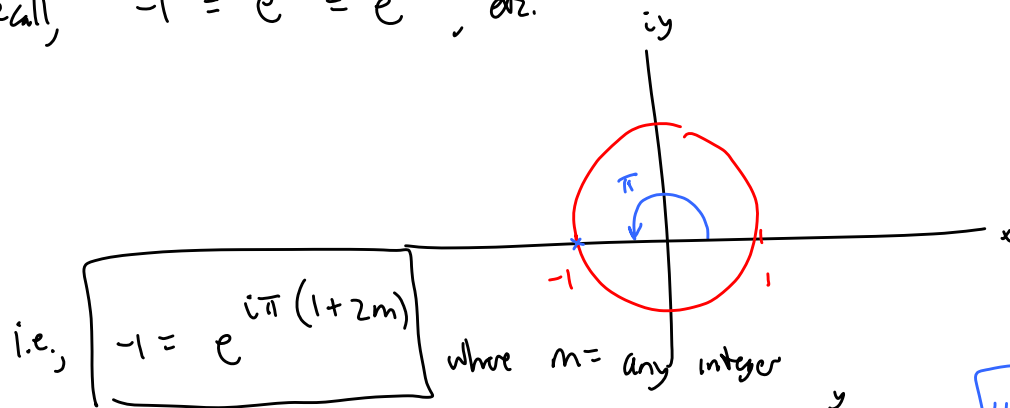
or \rightarrow if any roots are repeated, e.g., if q_1 is repeated

$$W(x) = C_1 e^{q_1 x} + C_2 x e^{q_1 x} + \dots$$

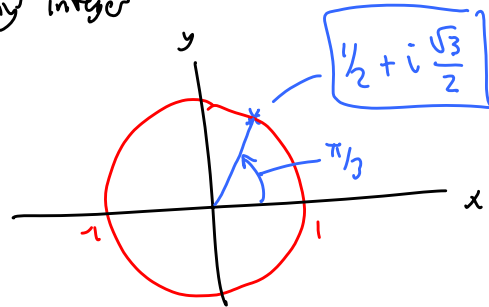
Here, we get 6 distinct roots $\Rightarrow (q^2 - K^2)^3 = -R_a K^2$

$$q = \pm K \sqrt[3]{1 + \left(\frac{R_a}{K^4}\right)^{1/3} (-1)^{1/3}} \quad (13 r)$$

Recall, $-1 = e^{i\pi} = e^{3i\pi}$, etc.

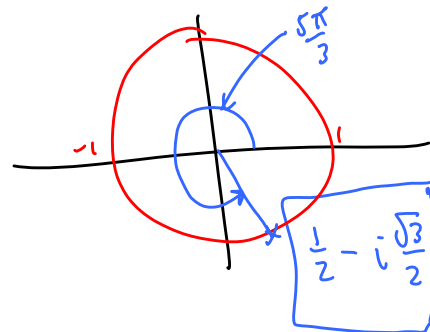


When $m=0$ $\therefore (-1)^{1/3} = (e^{i\pi})^{1/3} = e^{i\pi/3}$



When $m=1$ $(-1)^{1/3} = (e^{3i\pi})^{1/3} = e^{i\pi} = -1$

When $m=2$ $(-1)^{1/3} = (e^{5i\pi})^{1/3} = e^{5i\pi/3}$



\therefore 6 distinct \wedge roots of Eq. 13 r independent

Kundu calls them $\pm iq_0$, $\pm q$, & $\pm q^*$

Where

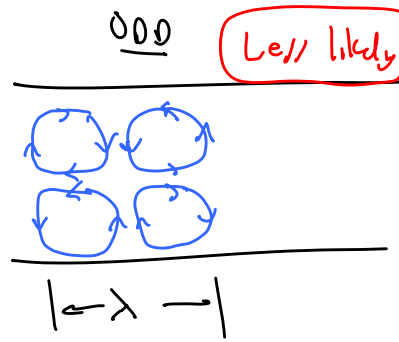
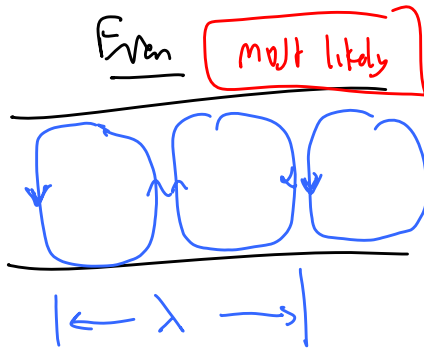
$$q_0 = K \sqrt{\left(\frac{R_0}{K^4}\right)^{1/3} - 1}$$

$$q = K \sqrt{1 + \frac{1}{2}\left(\frac{R_0}{K^4}\right)^{1/3} (1 + i\sqrt{3})}$$

$$q^* = \dots \dots \dots - \dots$$

So finally
$$W(z) = \underbrace{C_1 e^{iq_0 z} + C_2 e^{-iq_0 z}}_{C_1 = C_2 \text{ to be even}} + \underbrace{C_3 e^{qz} + C_4 e^{-qz}}_{C_3 \text{ must} = C_4 \text{ to be even}} + \underbrace{C_5 e^{q^* z} + C_6 e^{-q^* z}}_{C_5 \text{ must} = C_6 \text{ to be even}}$$

Also, see Kundu \rightarrow odd mode is less likely than the even mode



So, combine into $\cos(q_0 z)$ & $\cosh(qz)$

$$W(z) = A \cos(q_0 z) + B \cosh(qz) + C \cosh(q^* z)$$

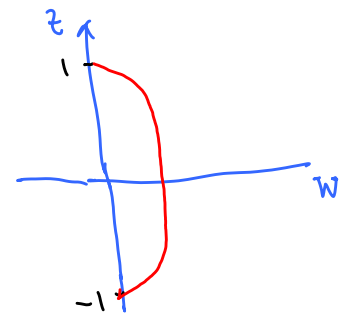
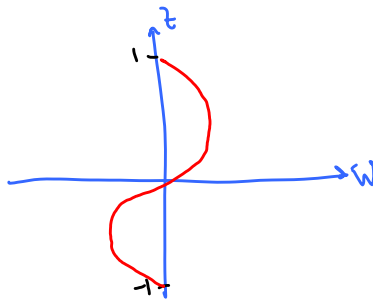
where A, B, C = new constants

A - even mode only

NOTE: AFTER CLASS THERE WAS A QUESTION ABOUT HOW WE ARE DEFINING ODD & EVEN HERE.

In terms of $W(z)$, the even mode has

The odd mode has



[As can be verified on the cell patterns above]